

Solutions - Hw #5

1) SCHROEDER 2.38

Schroeder's description of an ideal

mixture suggests the existence of

a multiplying function $\mathcal{R}(N)$ which

describes the multiplicity of a mixture

of N_A molecules of species A and N_B

molecules of species B by the rule:

$$\mathcal{R}_{A \otimes B}^{\text{DIST}}(N_A, N_B) = \mathcal{R}(N_A + N_B)$$

This multiplicity function $\mathcal{R}_{A \otimes B}$ is valid

for the Purely Hypothetical case where

different molecules of the same species

are distinguishable. Since the two

species have identical physical properties

except for their identity, it makes

Sense that the multiplicity of their mixture would then depend only on

these total particles # $N_A + N_B$.

To modify $\sqrt{N_{A+B}}$ to account for

the indistinguishability of molecules of

the same species we divide by the

number of permutations mixing molecules

of the same species, i.e. $N_A!$ for

the N_A molecules of species A and

N_B for species B. In other words:

$$\text{in dist} \quad \sqrt{N_{A+B}}(N_A, N_B) = \frac{\sqrt{N_A! N_B!}}{(N_A + N_B)}$$

$$= \frac{\sqrt{(N_A + N_B)!}}{N_A! N_B!}$$

The entropy of mixing is the difference in entropy between a mixture of A and B and a pure system containing only A or B, i.e.

$$\Delta S_{\text{mix}} = S(N_A, N_B) - S(N_A + N_B, 0)$$

$$= k_B \ln \mathcal{V}_{A+B}^{\text{INDIST}}(N_A, N_B)$$

$$- k_B \ln \mathcal{V}_{A+B}^{\text{INDIST}}(N_A + N_B, 0)$$

$$\text{expand } k_B \ln \left[\frac{\mathcal{V}(N_A + N_B)}{N_A! N_B!} \right]$$

$$- k_B \ln \left[\frac{\mathcal{V}(N_A + N_B)}{(N_A + N_B)! 0!} \right]$$

$$= k_B \ln \left[\frac{N!}{N_A! (N - N_A)!} \right] \quad N = N_A + N_B$$

$$= k_B \ln \left(\frac{N}{N_A} \right)$$

from Stirling's Approximation we
have: $\ln \left(\frac{n}{N_A} \right)$

$$= \ln \frac{n!}{N_A! (n-N_A)!}$$

$$= \ln n! - \ln N_A! - \ln (n-N_A)!$$

$$= N \ln N - N - N_A \ln N_A + N_A$$

$$- (n-N_A) \ln (n-N_A) + N - N_A$$

$$= N \left[\ln N - x \left[\ln N + \ln x \right] \right. \\ \left. - (1-x) \left[\ln N + \ln (1-x) \right] \right]$$

$$= -N \left[x \ln x + (1-x) \ln (1-x) \right]$$

$$x = \frac{N_A}{n}$$

2) a) Since the wall separating

A + B is fixed we have

$$dV_1 = 0, dV_2 = 0$$

Since the wall is closed [no holes]

we have

$$dN_1 = 0, dN_2 = 0$$

From energy conservation we have

$$du_1 = -du_2 \equiv du$$

so that:

$$dS_{\text{TOT}} = \sum_{j=1}^2 \frac{\partial S_{\text{TOT}}}{\partial u_j} du_j + \frac{\partial S}{\partial V_j} dV_j + \frac{\partial S}{\partial N_j} dN_j$$

$$= \left[\frac{\partial S_{\text{TOT}}}{\partial u_1} - \frac{\partial S_{\text{TOT}}}{\partial u_2} \right] du$$

$$= \left[\frac{\partial S_1}{\partial u_1} - \frac{\partial S_2}{\partial u_2} \right] du$$

$$= \left[\frac{1}{T_1} - \frac{1}{T_2} \right] du$$

So that the energy transfer is

irreversible unless T_1 is very close

to T_2 [if T_1 is exactly T_2 then

no energy transfer occurs at all].

^{1 AGAIN}
b) Here we have a closed wall so

$$dN_1 = 0, \quad dN_2 = 0$$

The wall can move but the total
volume of the combined system

is fixed so

$$dV_1 = -dV_2 \equiv dV$$

Again Energy is conserved so

$$dU_1 = -dU_2 \equiv dU$$

We also assume the thermal contact

is good so that

$$dT_1 = dT_2 \quad \text{or}$$

$$\frac{\partial S_1}{\partial U_1} dU_1 + \frac{\partial S_2}{\partial U_2} dU_2 = 0$$

$$\frac{\partial S_1}{\partial U_2} \frac{\partial S_2}{\partial U_2} dU_1 dU_2 = 0$$

~~After 3 weeks~~

$$d \left[\frac{u_1}{\frac{3}{2} N_1 k_B} \right] = d \left[\frac{u_2}{\frac{3}{2} N_2 k_B} \right]$$

$$\frac{1}{N_1} du_1 = \frac{1}{N_2} du_2$$

$$\frac{du}{N_1} = -\frac{du}{N_2}$$

$$du \left(\frac{1}{N_1} + \frac{1}{N_2} \right) = 0$$

$$du = 0$$

so :

$$dS_{TOT} = \left(\frac{\partial S_1}{\partial V_1} - \frac{\partial S_2}{\partial V_2} \right) dV$$

≈ But ~~per~~ from Sackur Tetrode Eqⁿ :

$$\frac{\partial}{\partial V} S = \frac{\partial}{\partial V} Nk_B \ln V + f(u, n)$$

≈ no V
dep.

$$= \frac{Nk_B}{V} = \frac{P}{T}$$

$$\text{so } dS_{\text{tot}} = \left(\frac{P_1}{T_1} - \frac{P_2}{T_2} \right) dv$$

But $T_1 = T_2 \equiv T$ so

$$dS_{\text{tot}} = \frac{1}{T} \int [P_1 - P_2] dv$$

so process is irreversible unless

P_1 is very close to P_2 [if P_1 is exactly P_2 no volume transfer occurs].

c) Now we have:

$$du_1 = -du_2 \equiv du \quad [\text{energy cons.}]$$

$$dV_1 = dV_2 = 0 \quad [\text{fixed vol.}]$$

$$dN_1 = -dN_2 \equiv dN \quad [\text{particle \# cons.}]$$

$$dT_1 = dT_2$$

~~$$\Rightarrow d\left(\frac{3}{2}N_1k_B T_1\right)$$~~

~~$$\text{But } d\left(\frac{u_1}{\frac{3}{2}N_1k_B}\right) = d\left(\frac{u_2}{\frac{3}{2}N_2k_B}\right)$$~~

$$\frac{du_1}{N_1} - \frac{u_1}{N_1^2} dN_1 = \frac{du_2}{N_2} - \frac{u_2}{N_2^2} dN_2$$

~~$$\text{But } \frac{du}{N} \sim \frac{du}{N}$$~~

$$= \frac{du}{N} - \frac{3}{2} kT \frac{dN}{N}$$

$$= \frac{3}{2} kT \left(\frac{du}{N} - \frac{dN}{N} \right)$$

$$\text{so } dT_1 = dT_2 \Rightarrow du \left[\frac{1}{u_1} + \frac{1}{u_2} \right] = dN \left[\frac{1}{N_1} + \frac{1}{N_2} \right]$$

so since

$$dS = \frac{\partial S}{\partial U} dU + \frac{\partial S}{\partial N} dN$$

and $\frac{\partial S}{\partial U} = \frac{3}{2} \frac{Nk_B}{U} = \frac{3}{2} \frac{Nk_B}{Nk_B T \cdot \frac{3}{2}} = \frac{1}{T}$

and $\frac{\partial S}{\partial N} = -\frac{5}{2} \frac{Nk_B}{n} + \frac{s}{n}$

$$= \frac{s}{n} - \frac{5}{2} k_B$$

we get

$$dS_{tot} = \frac{1}{T} dU_1 + \left[\frac{s_1}{n_1} - \frac{5}{2} k_B \right] dN_1$$

$$+ \frac{1}{T} dU_2 + \left[\frac{s_2}{n_2} - \frac{5}{2} k_B \right] dN_2$$

$$\sim dU_1 = -dU_2$$

$$= \left[\frac{s_1}{n_1} - \frac{s_2}{n_2} \right] dN$$

so in this case we find the Molar Entropies need to be nearly equal for the process to be reversible.

NOTE: It turns out that in general

$$\mu - \frac{\partial S}{\partial N} = \text{const}$$

it is the $\frac{\partial S}{\partial N}$ which must be equal

for reversible particle transfer.

There is a quantity

$$\mu = \cancel{-T} - T \frac{\partial S}{\partial N}$$

called the chemical potential which

also serves to determine the

irreversibility of a ~~process~~ particle

transfer.